

# Long-range spin transport in superconductors

Detlef Beckmann

Florian Hübler, Michael J. Wolf

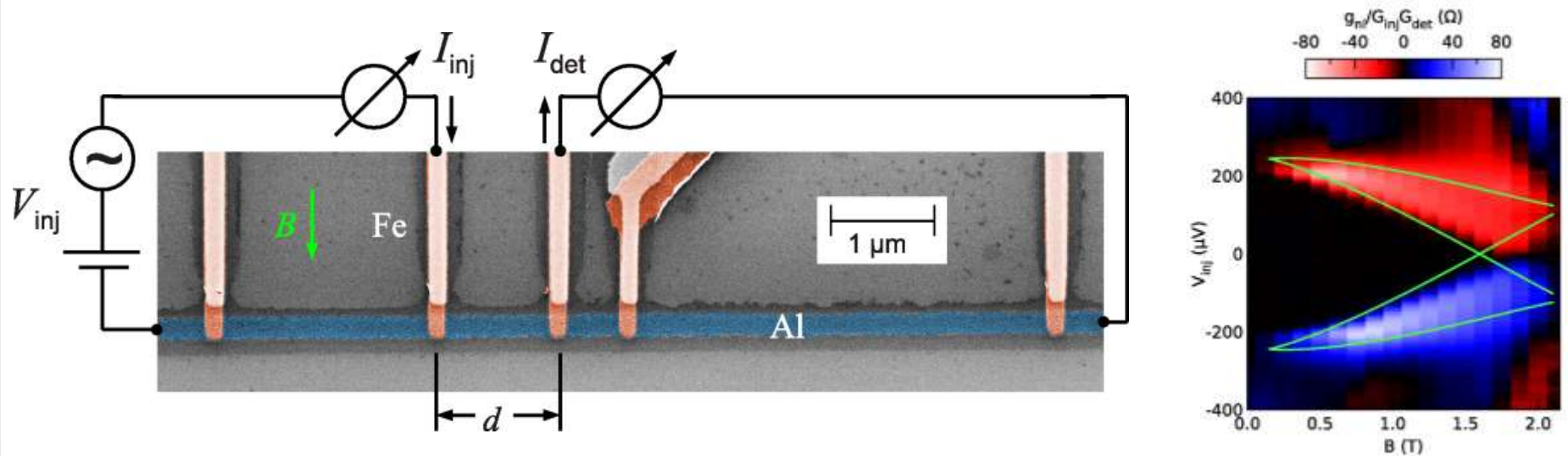
Julien Camyrand Lemyre

Nikolai Gulnizkij

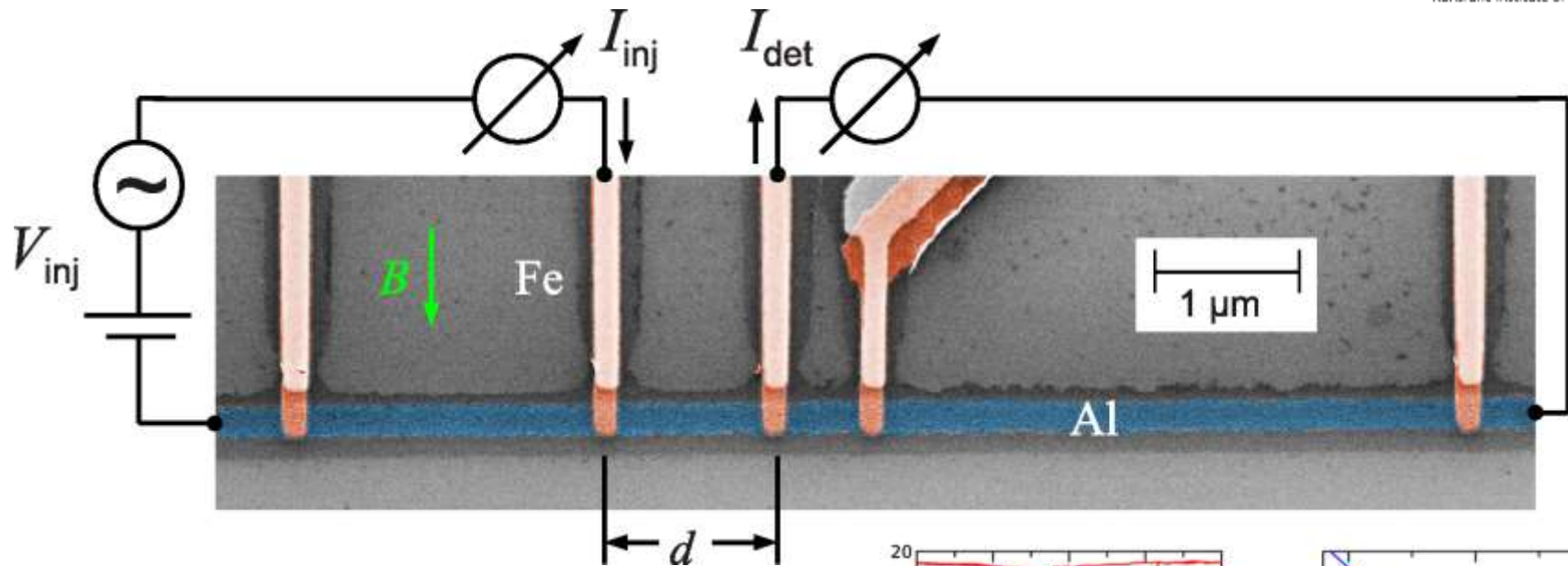
Stefan Kolenda

Hilbert v. Löhneysen

INSTITUTE OF NANOTECHNOLOGY

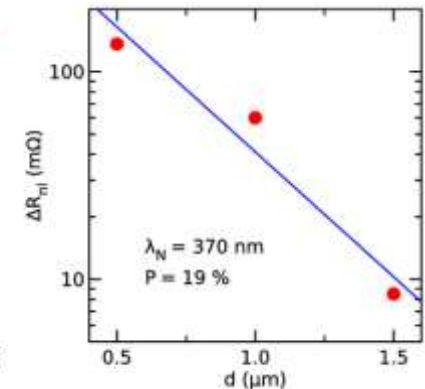
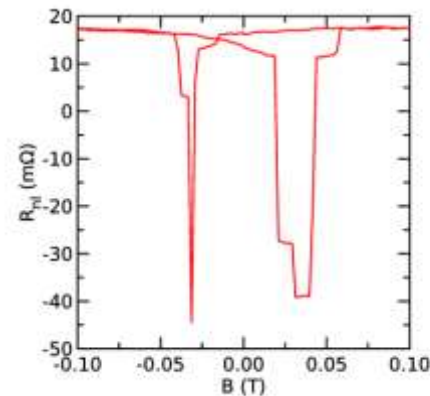


# Experiment – samples



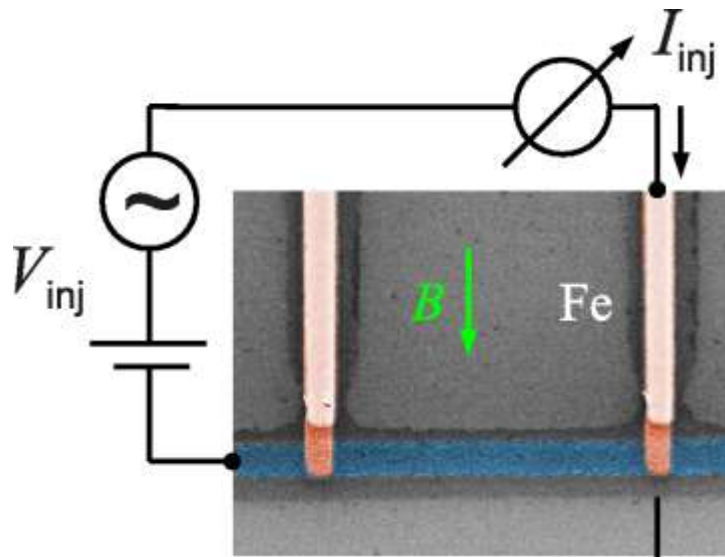
- film thickness  $\sim 12$  nm
- contact distances  $d \sim 0.5 \mu\text{m} - 8 \mu\text{m}$
- magnetic field  $B \parallel$  iron wires
- parallel magnetization alignment

Hübler et al., Phys. Rev. Lett. **109**, 207001 (2012)



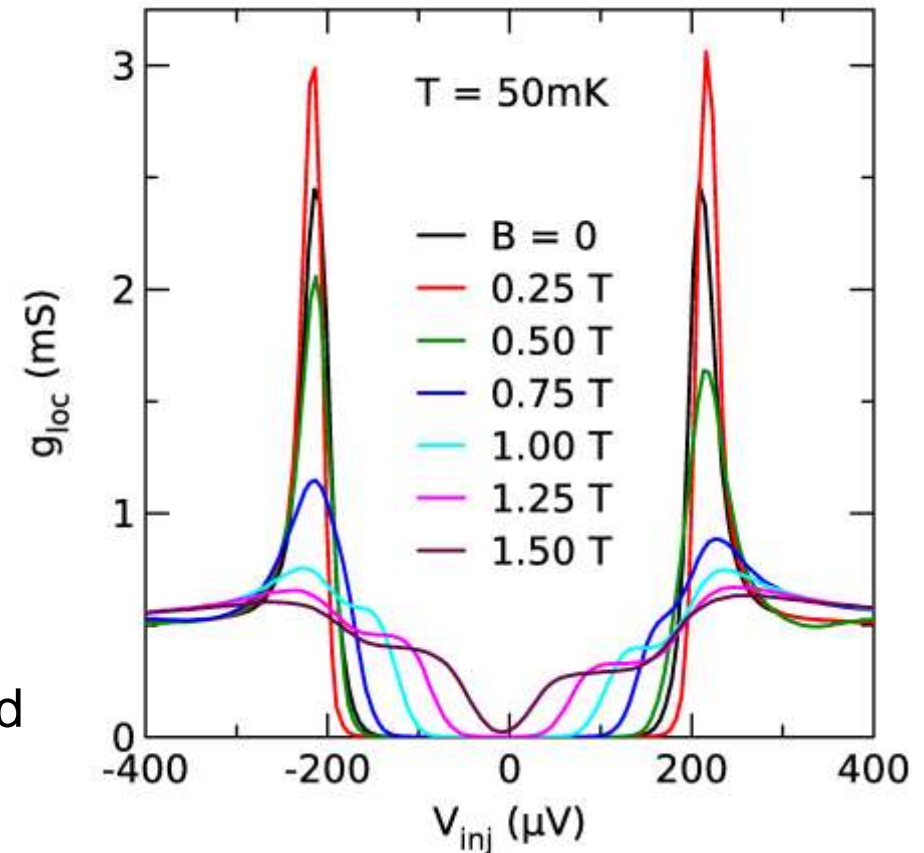
normal-state spin-valve:  
 $\rightarrow$  spin diffusion length  $\lambda_N = 370$  nm

# Experiment – *junction characterization*

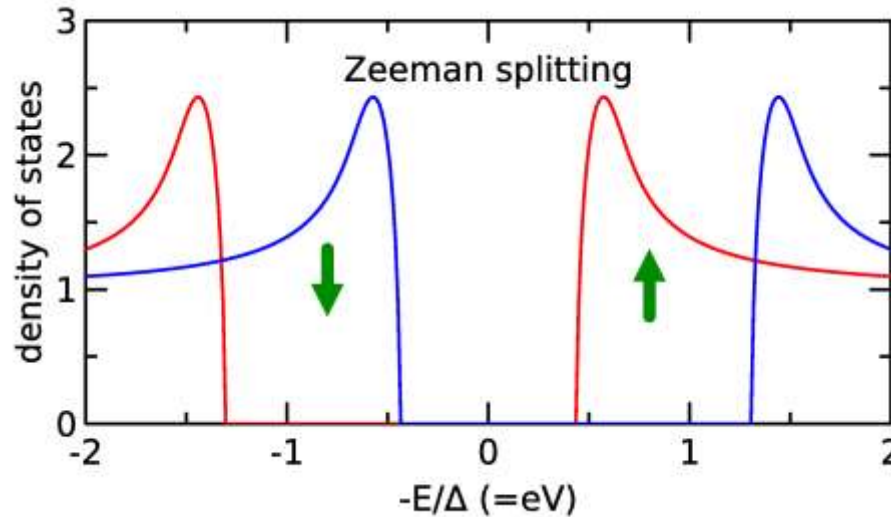


local conductance

Zeeman splitting in the magnetic field



# Model – Zeeman splitting



Zeeman-splitting of DOS (semiconductor model)

$$n_{\downarrow,\uparrow}(E) = n(E \pm \mu_B B)$$

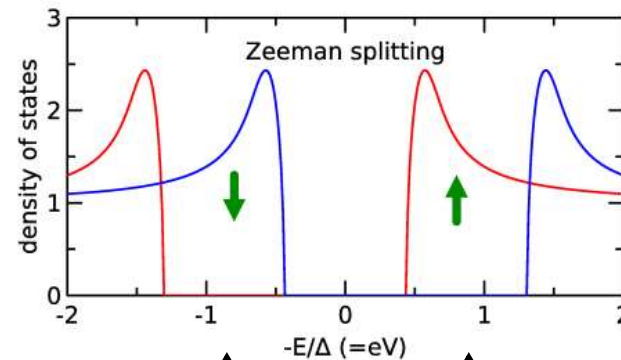
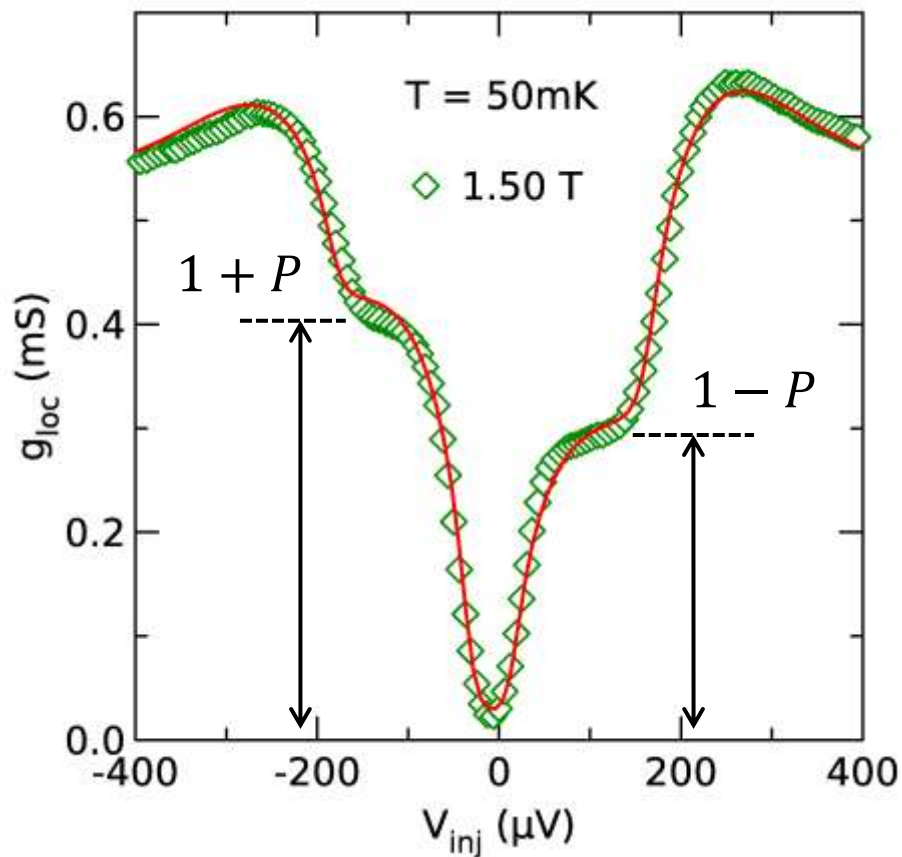
Charge current 
$$I = \frac{1}{e} \int [G_{\downarrow} n_{\downarrow}(E) + G_{\uparrow} n_{\uparrow}(E)] \{f(E - eV) - f(E)\} dE$$

Tedrow & Meservey, PRL **26**, 192 (1971)  
 Meservey & Tedrow, Phys. Rep. **238**, 173 (1994)

$$G_{\downarrow,\uparrow} = \frac{G_N}{2} (1 \pm P)$$

# Experiment – junction characterization

Al/Al<sub>2</sub>O<sub>3</sub>/Fe



superconductor acts as spin filter

asymmetry

→ determination of  $P = \frac{G_{\downarrow} - G_{\uparrow}}{G_{\downarrow} + G_{\uparrow}}$

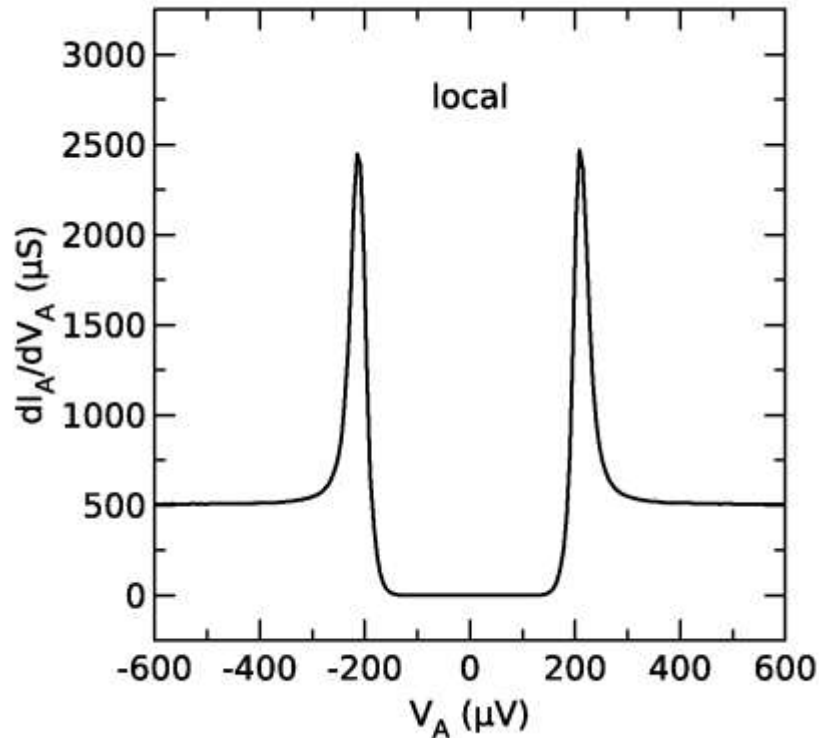
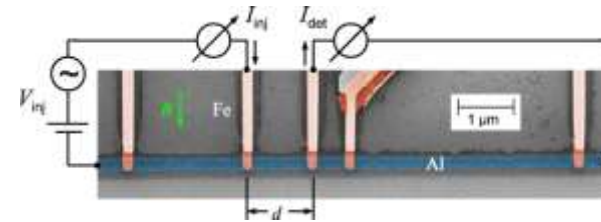
$P = 19\%$

(consistent with spin-valve experiments)

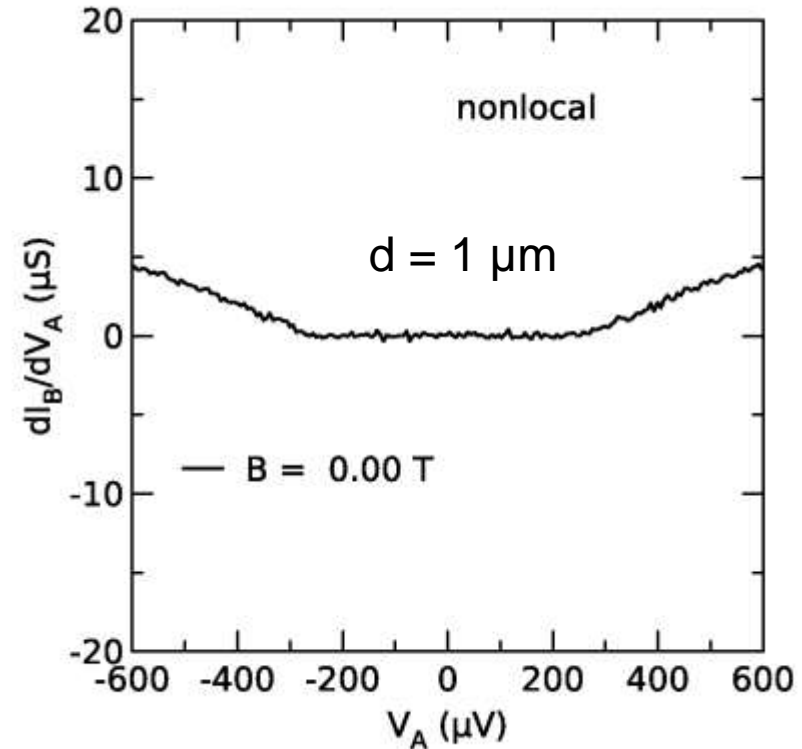
# Experiment – *local vs nonlocal conductance*

$B = 0$

Al/Al<sub>2</sub>O<sub>3</sub>/Fe



density of states

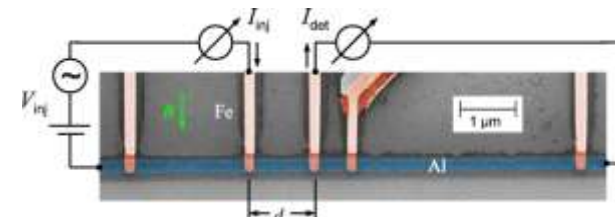


charge imbalance

Clarke, Phys. Rev. Lett. **28**, 1363 (1972)

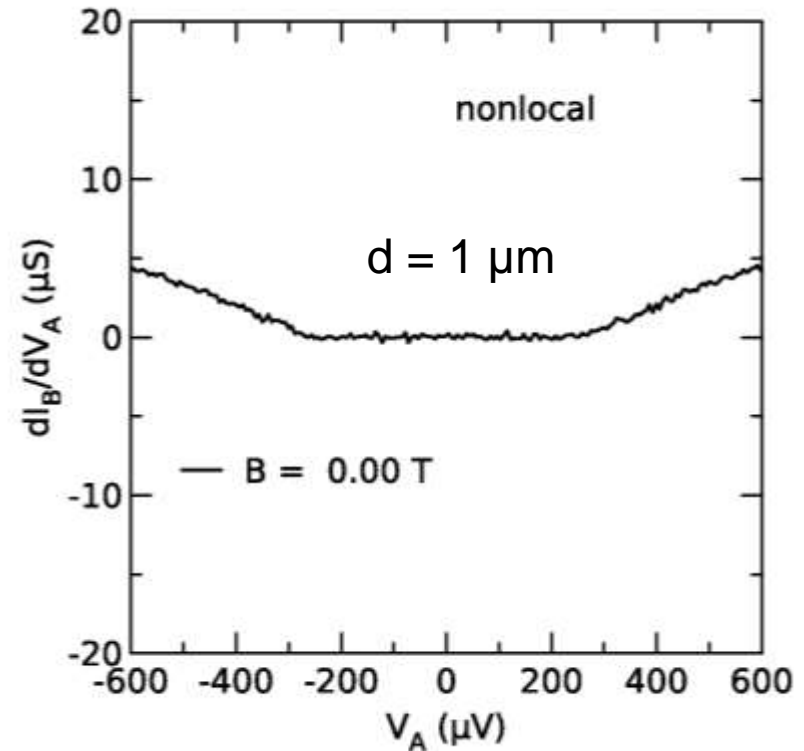
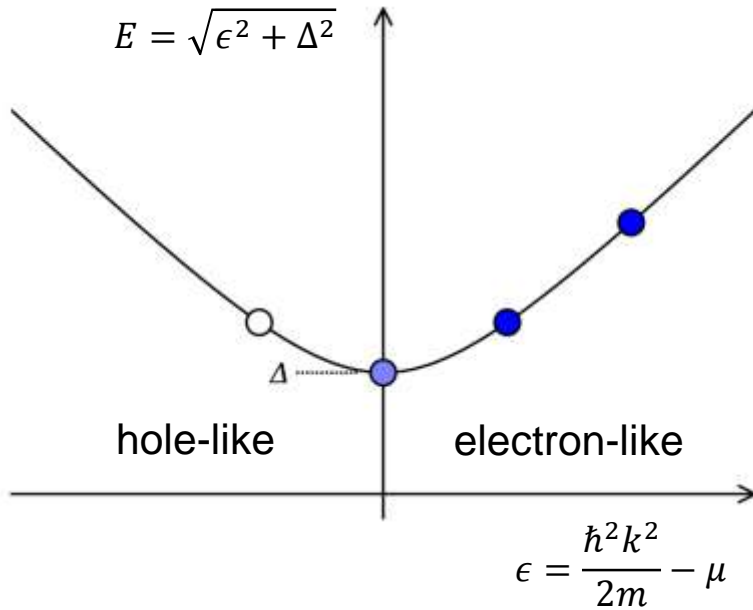
Tinkham & Clarke, Phys. Rev. Lett. **28**, 1366 (1972)

# Experiment – *local vs nonlocal conductance*



quasiparticle dispersion

$$E = \sqrt{\epsilon^2 + \Delta^2}$$



charge imbalance

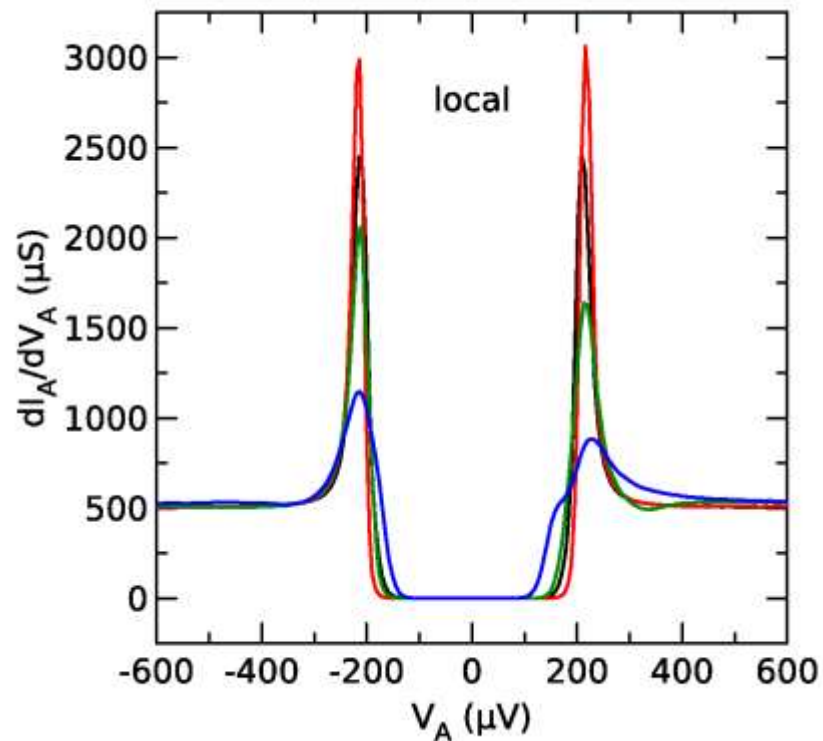
Clarke, Phys. Rev. Lett. **28**, 1363 (1972)

Tinkham & Clarke, Phys. Rev. Lett. **28**, 1366 (1972)

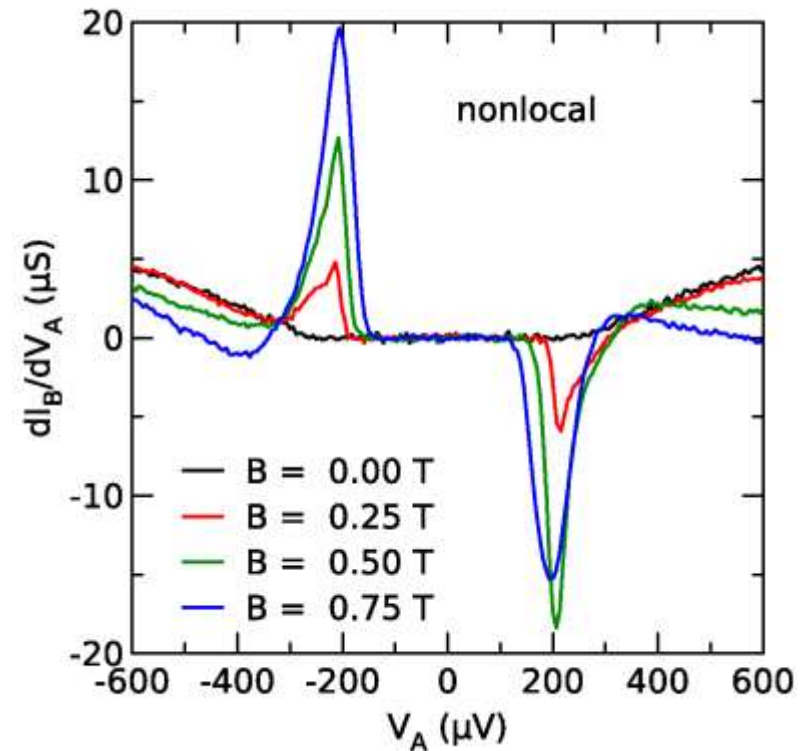
Hübler et al., Phys. Rev. B **81**, 184524 (2010)

# Experiment – *local vs nonlocal conductance*

$B > 0$



broadening & Zeeman splitting

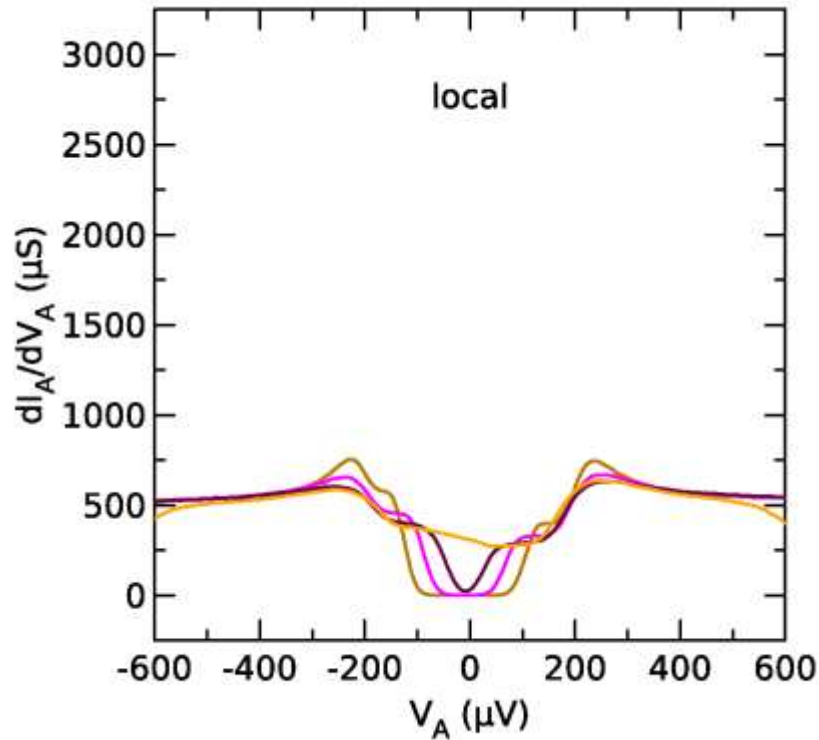


asymmetric features appear

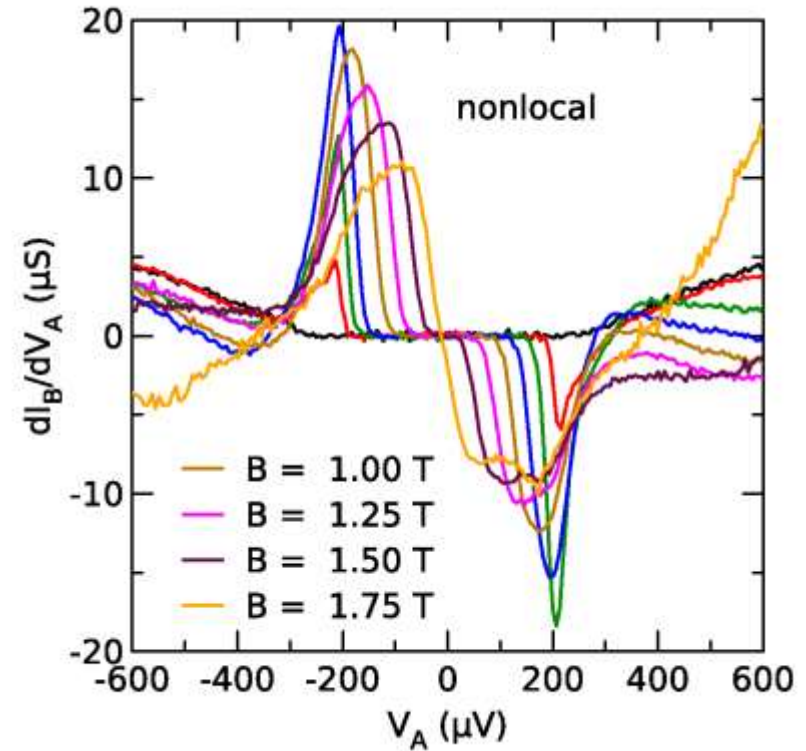


# Experiment – *local vs nonlocal conductance*

$B \gg 0$



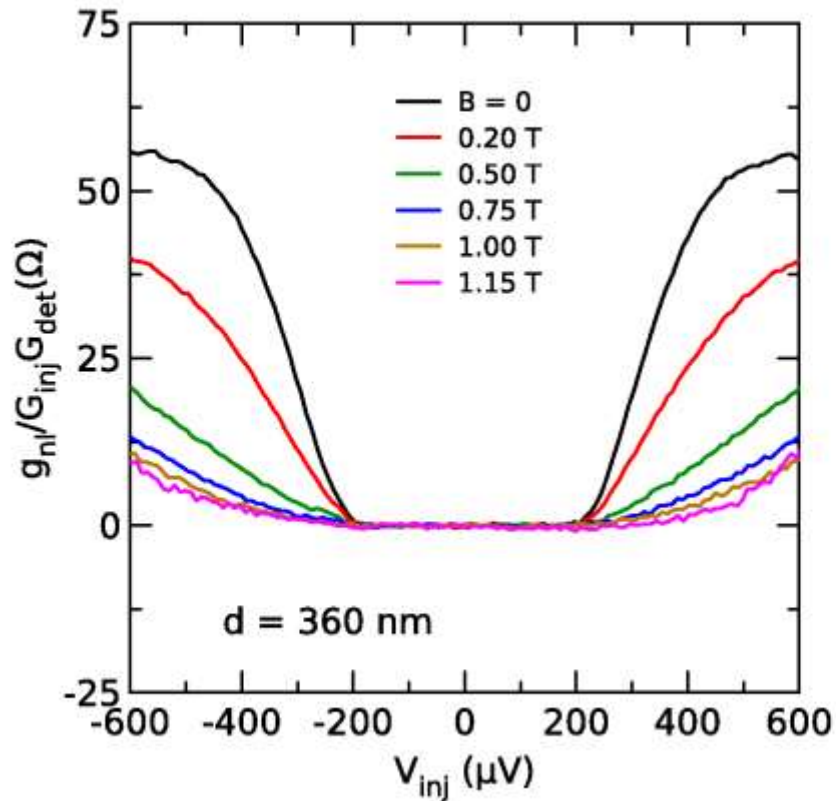
Zeeman splitting grows



features broaden

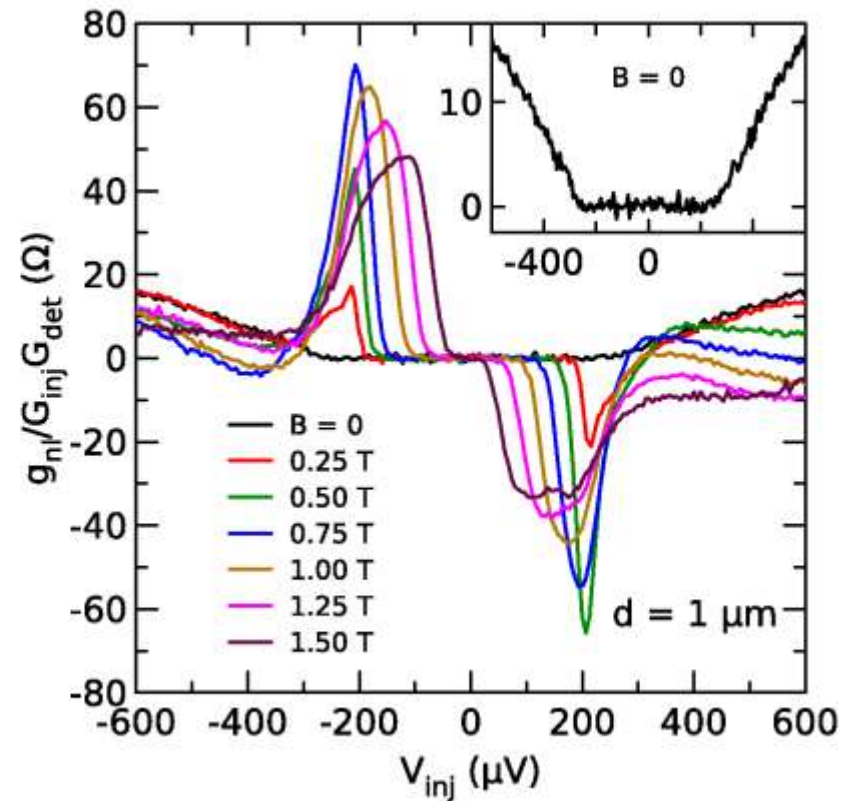
# Experiment – *normal vs ferromagnet*

## NISIN



only charge imbalance

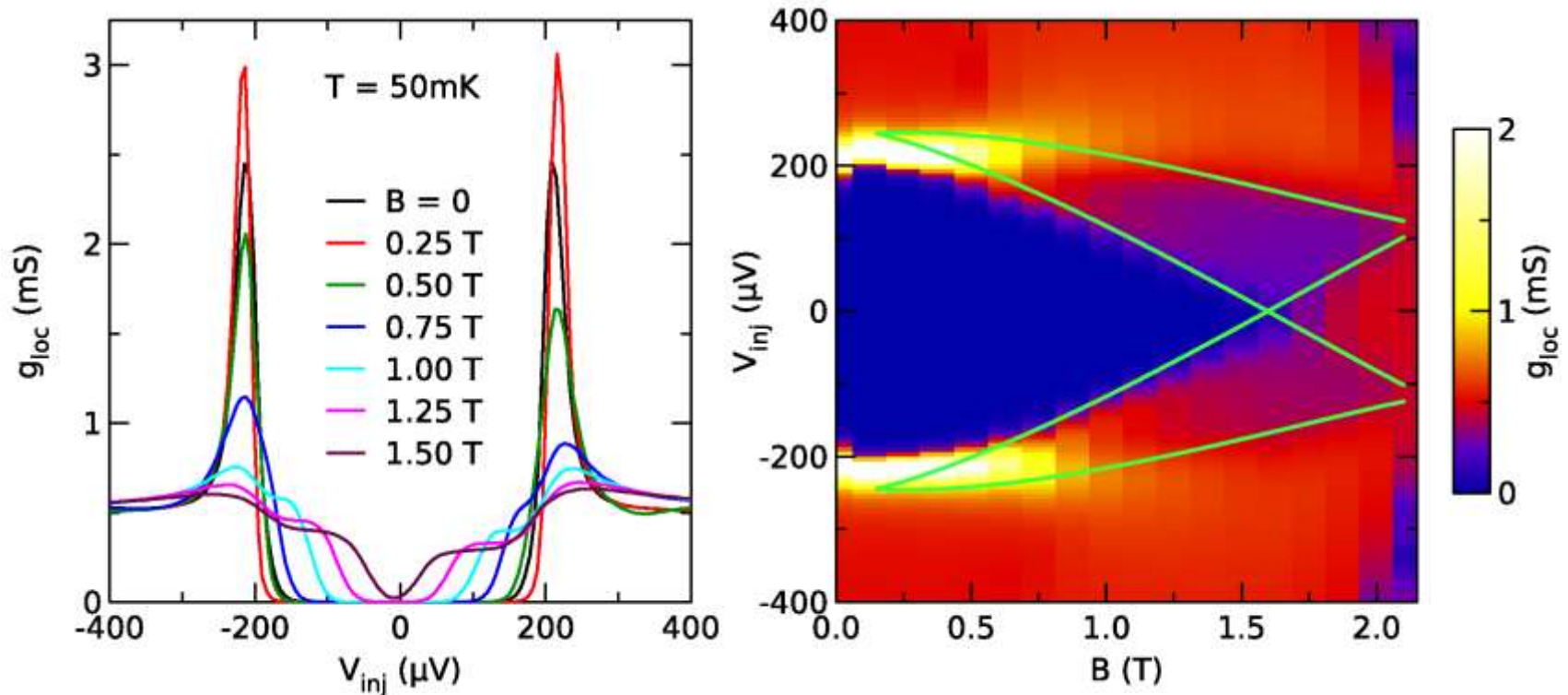
## FISIF



charge imbalance + asymmetry

similar results: Quay et al., Nat. Phys. **9**, 84 (2013)

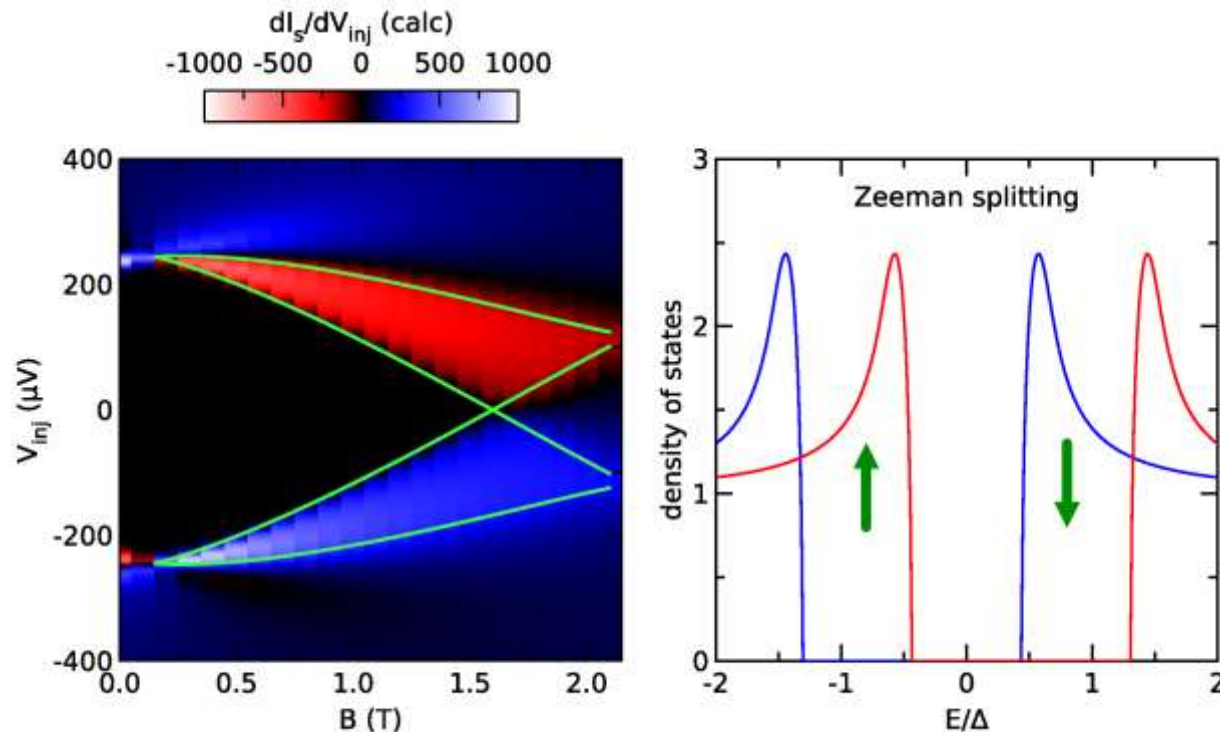
# Experiment – *local conductance*



$$I = \frac{1}{e} \int [G_{\downarrow} n_{\downarrow}(E) + G_{\uparrow} n_{\uparrow}(E)] \{f(E - eV) - f(E)\} dE$$

wedge-shaped regions: only one spin band contributes  
superconductor acts as spin filter

# Experiment – *predicted spin injection*

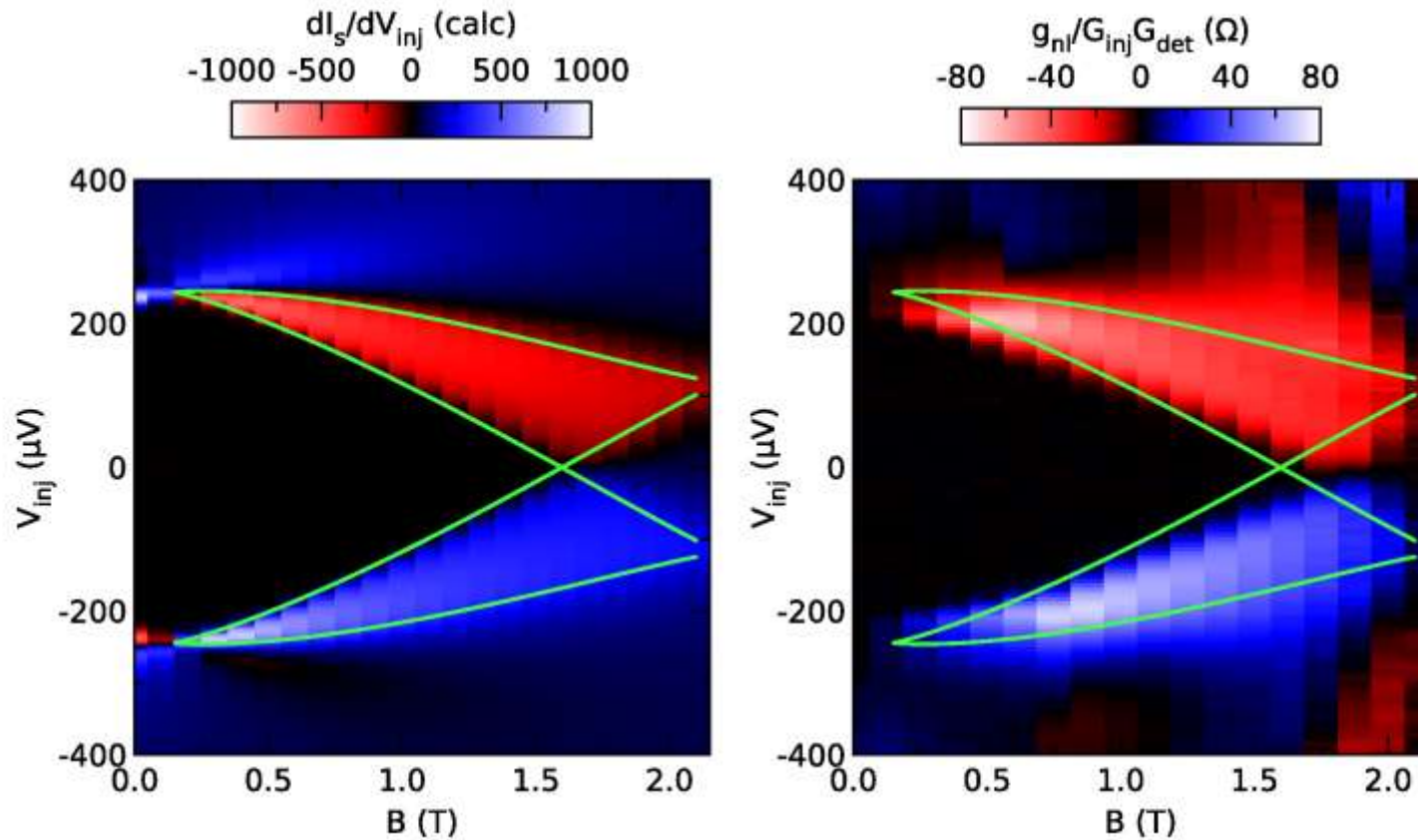


Spin current 
$$I_S = \frac{1}{e} \int [G_{\downarrow} n_{\downarrow}(E) - G_{\uparrow} n_{\uparrow}(E)] \{f(E - eV) - f(E)\} dE$$

$$I_S \propto \frac{G_{\downarrow} - G_{\uparrow}}{G_{\downarrow} + G_{\uparrow}} + \frac{n_{\downarrow} - n_{\uparrow}}{n_{\downarrow} + n_{\uparrow}}$$

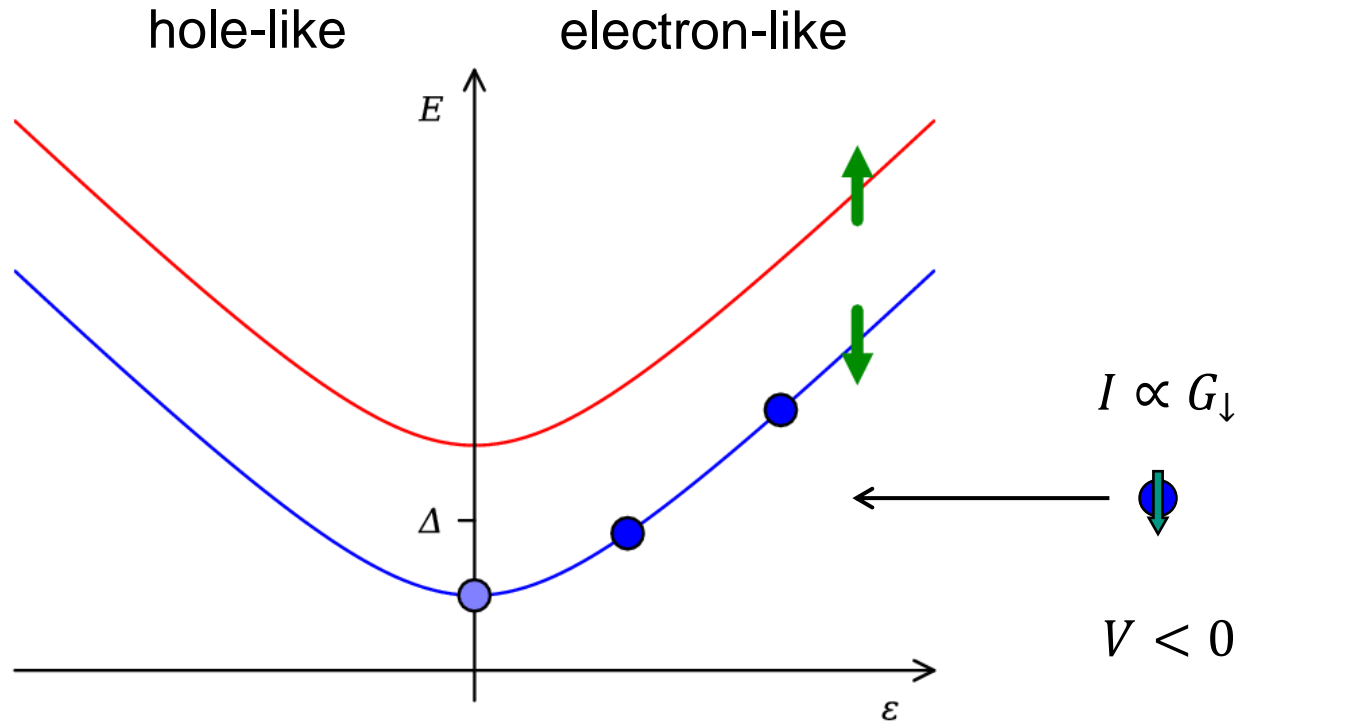
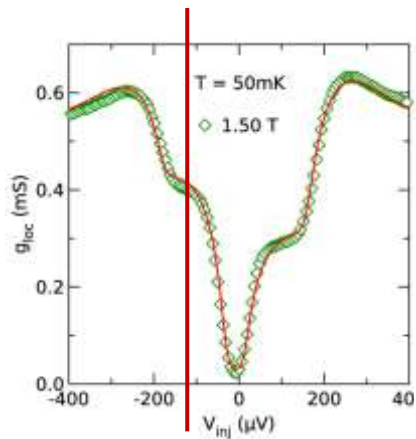
see also Giazotto & Taddei, Phys. Rev. B **77**, 132501 (2008)

# Experiment – *spin injection vs signal*



nonlocal conductance follows spin injection

# Model – injection

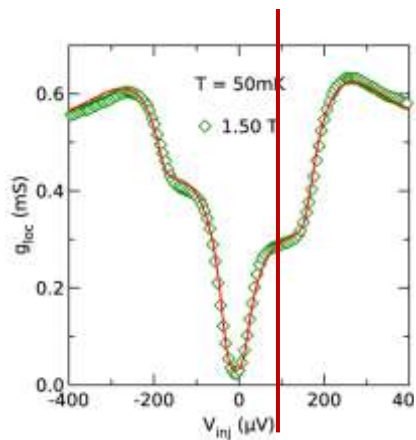


charge imbalance  $Q_{\downarrow}^* < 0$

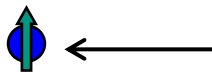
spin accumulation  $S_{\downarrow} > 0$

inject spin-down electrons

# Model – injection



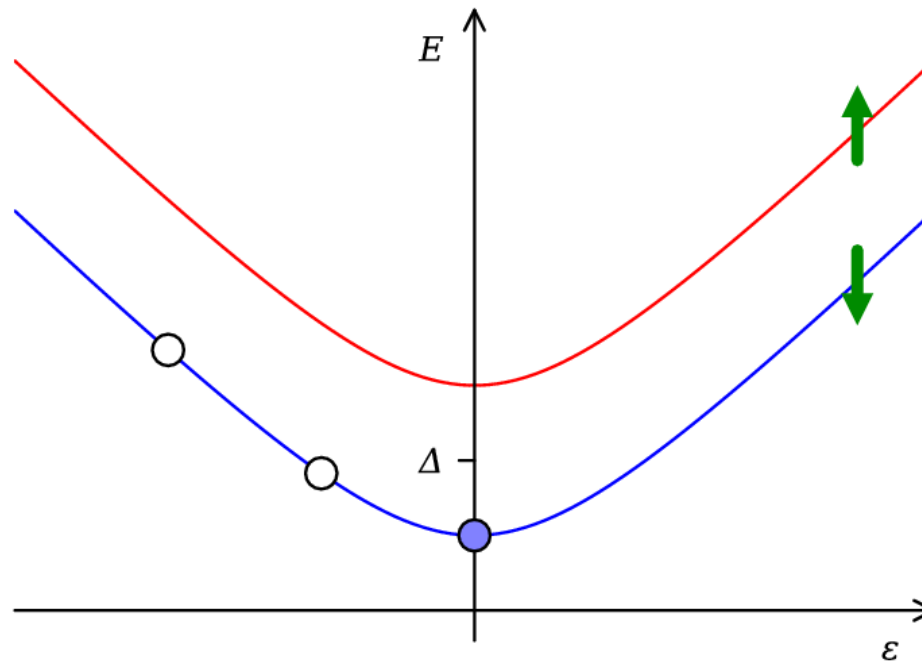
$$I \propto G_{\uparrow}$$



$$V > 0$$

hole-like

electron-like



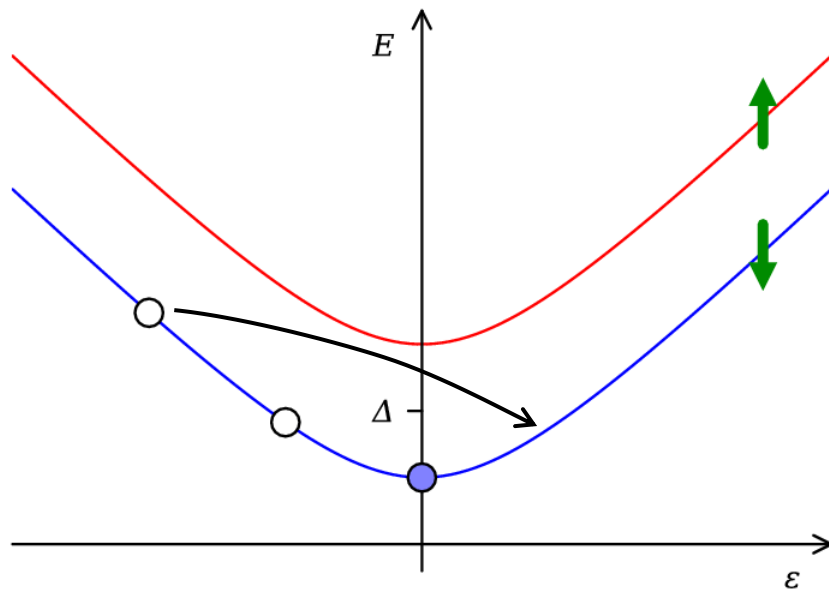
extract spin-up electrons

$$Q_{\downarrow}^* > 0$$

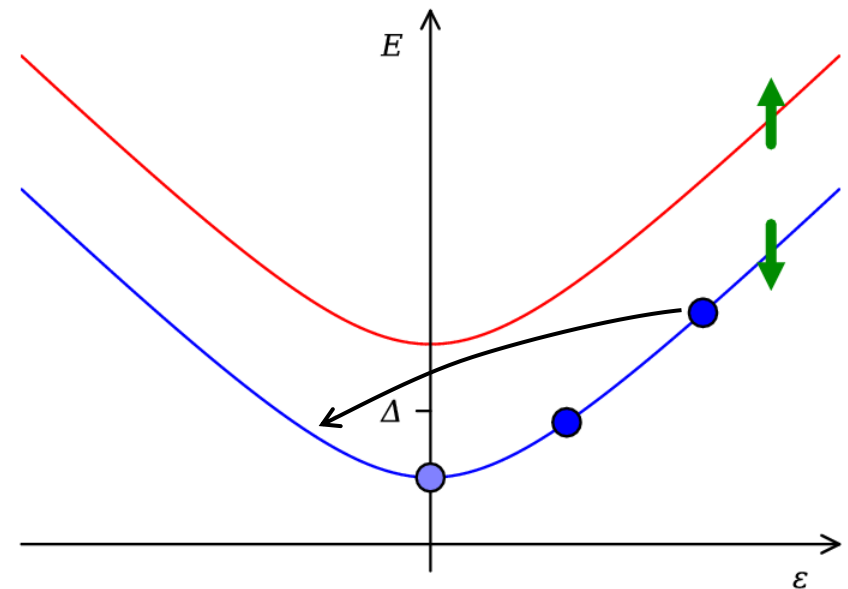
$$S_{\downarrow} > 0$$

# Model – charge relaxation

charge relaxation for  $d > \lambda_Q^*$



$V > 0$



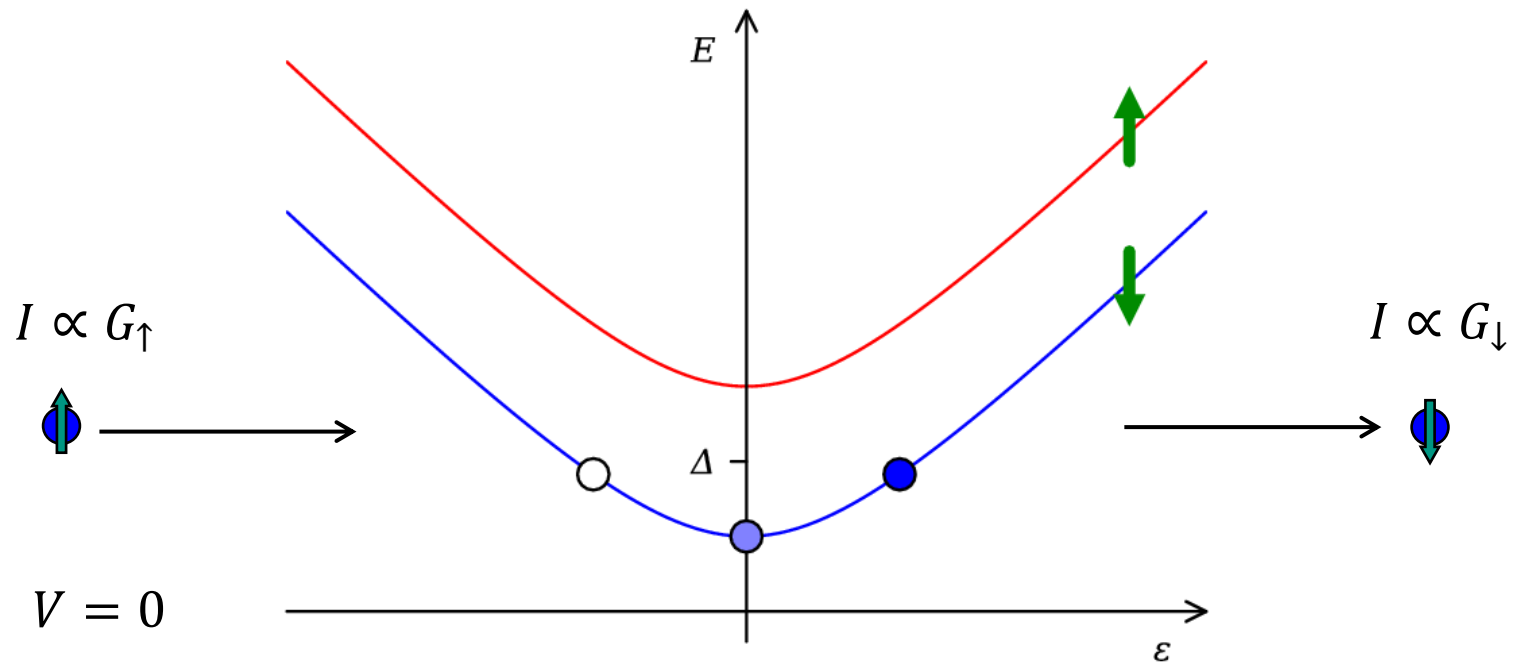
$V < 0$

fast for  $E \gtrsim \Delta$



# Model – detector current

$$I_{\text{det}} = (G_{\downarrow} + G_{\uparrow})(Q_{\downarrow}^* + Q_{\uparrow}^*) - (G_{\downarrow} - G_{\uparrow})(S_{\downarrow} - S_{\uparrow}) \approx -(G_{\downarrow} - G_{\uparrow})S_{\downarrow}$$



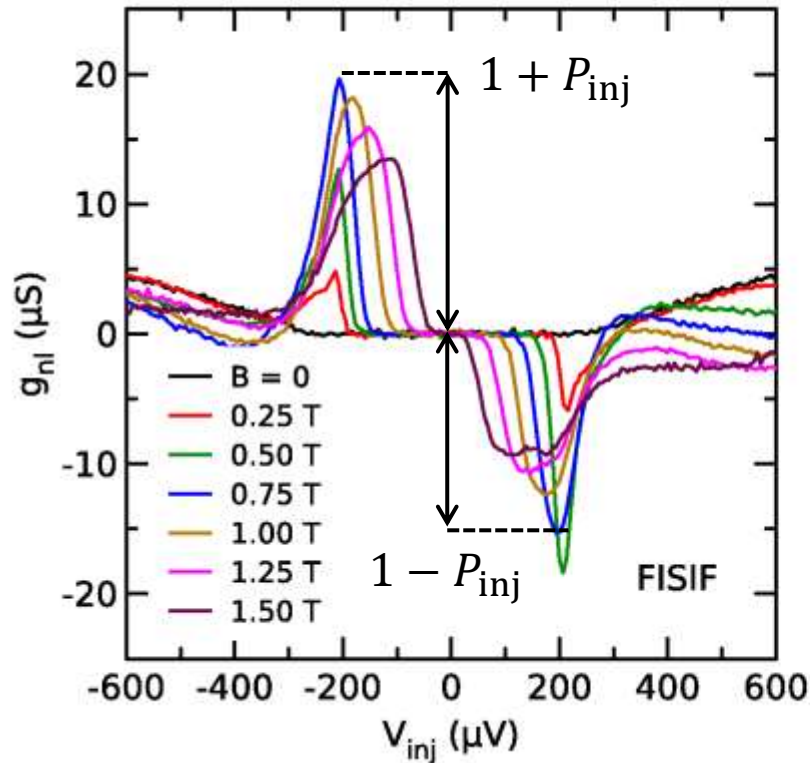
$$Q_{\downarrow}^* \approx 0$$

$$S_{\downarrow} > 0$$

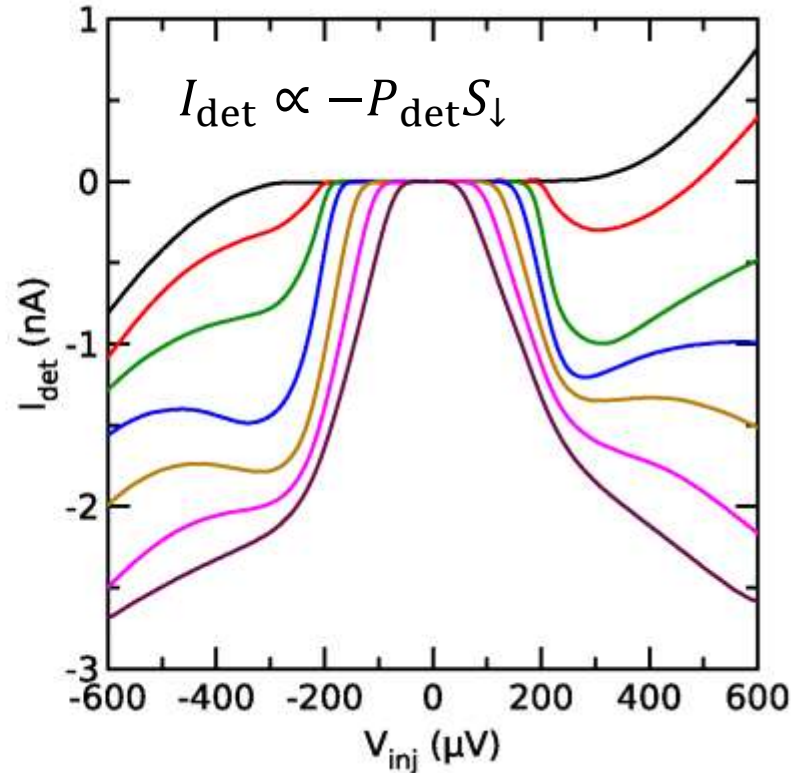
spin charge separation

see Zhao & Hershfield, PRB **52**, 3632 (1995)

# Interpretation – *detector current*

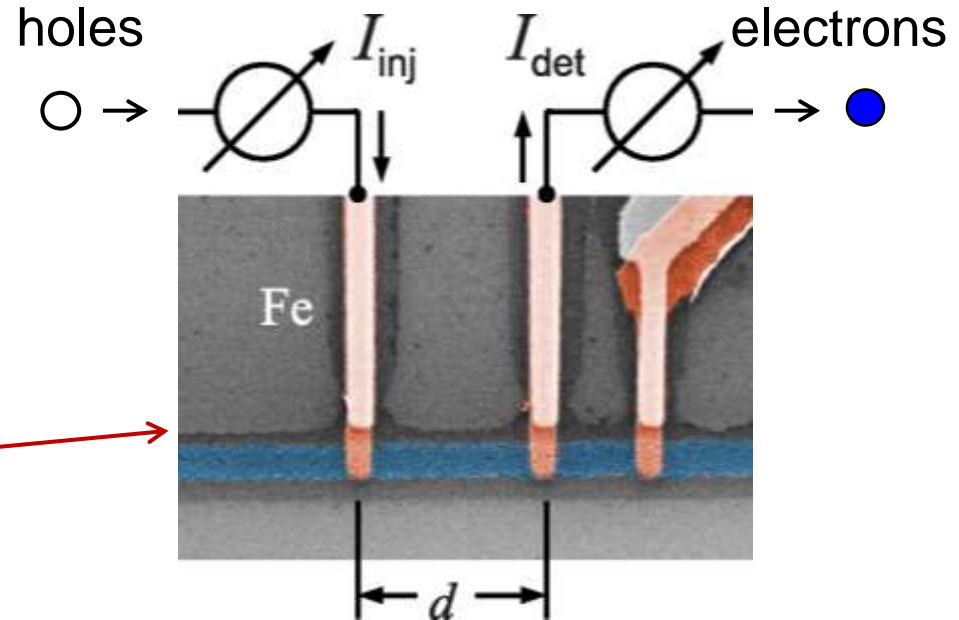
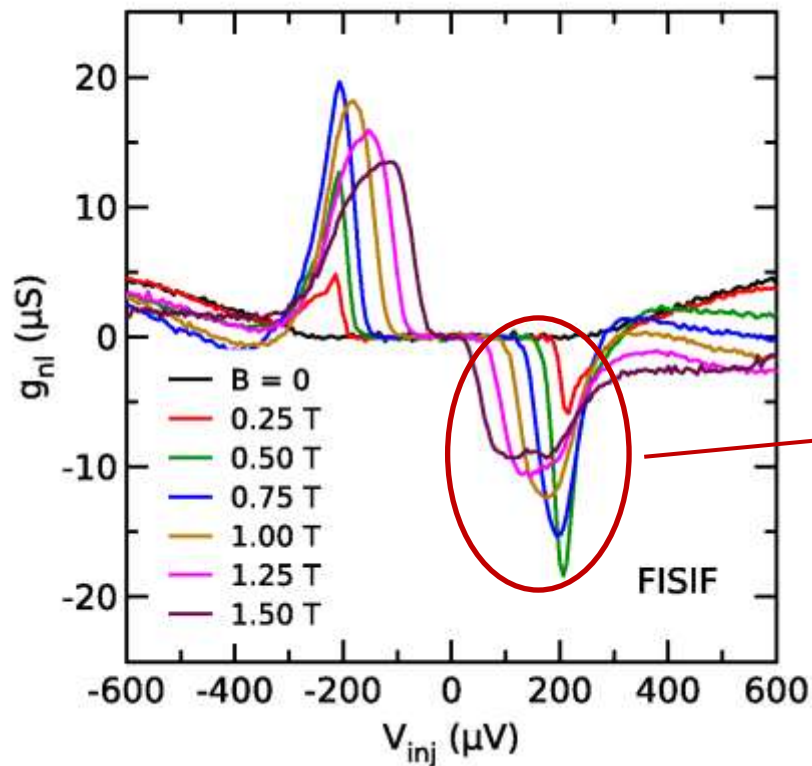


injector polarisation  $P_{inj}$   
 → different peak heights

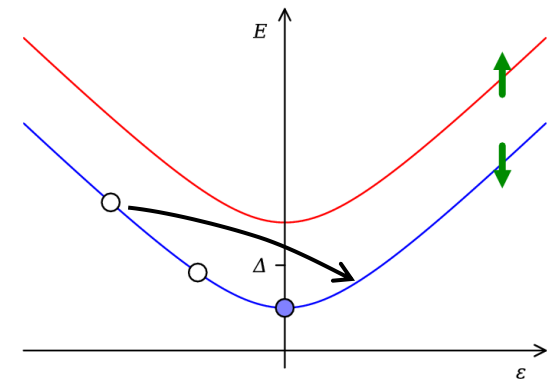


detector polarisation  $P_{det}$   
 → sign of nonlocal current

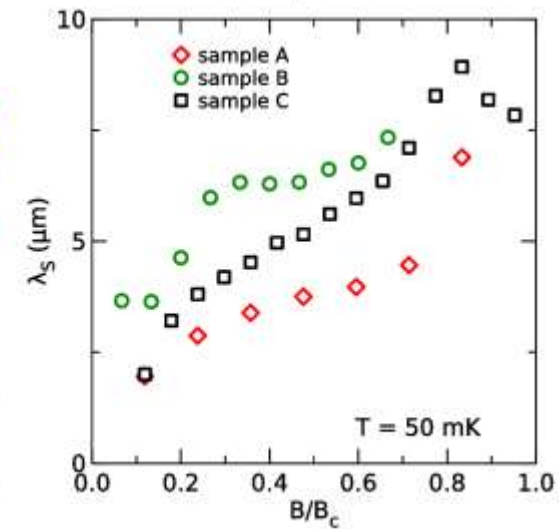
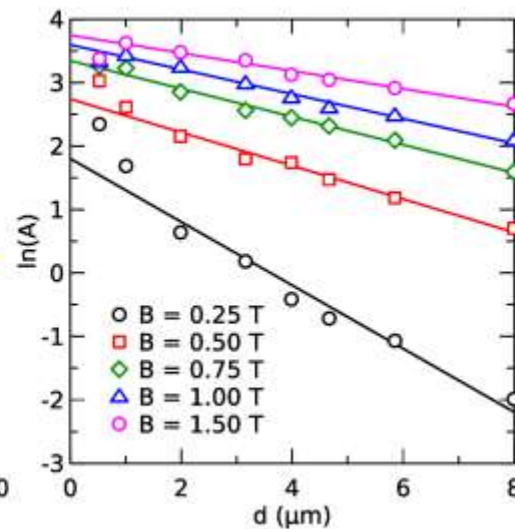
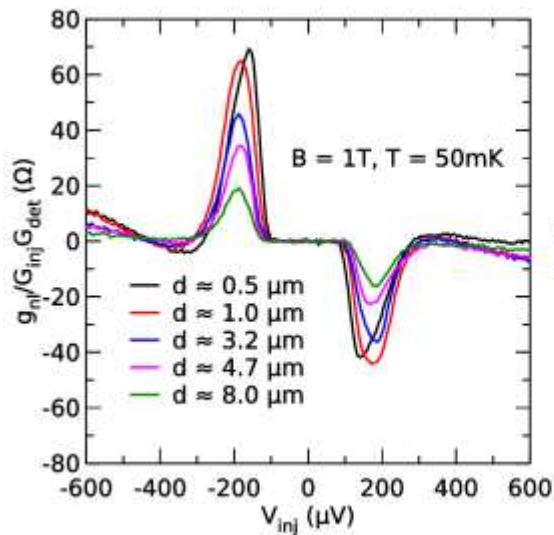
# Interpretation – *negative nonlocal signal*



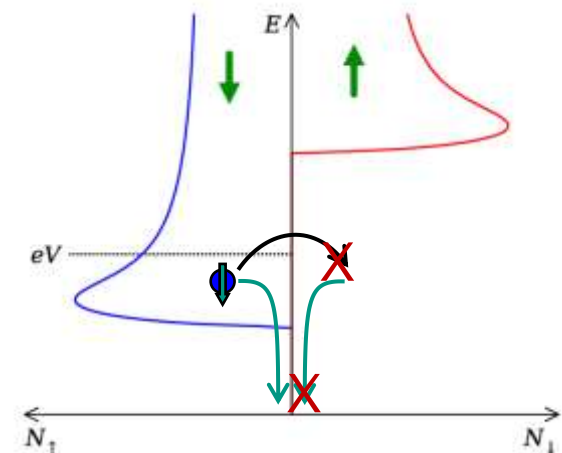
at positive bias: inject holes, get electrons.  
 looks like crossed Andreev reflection. why?  
 charge relaxation (no coherence needed)



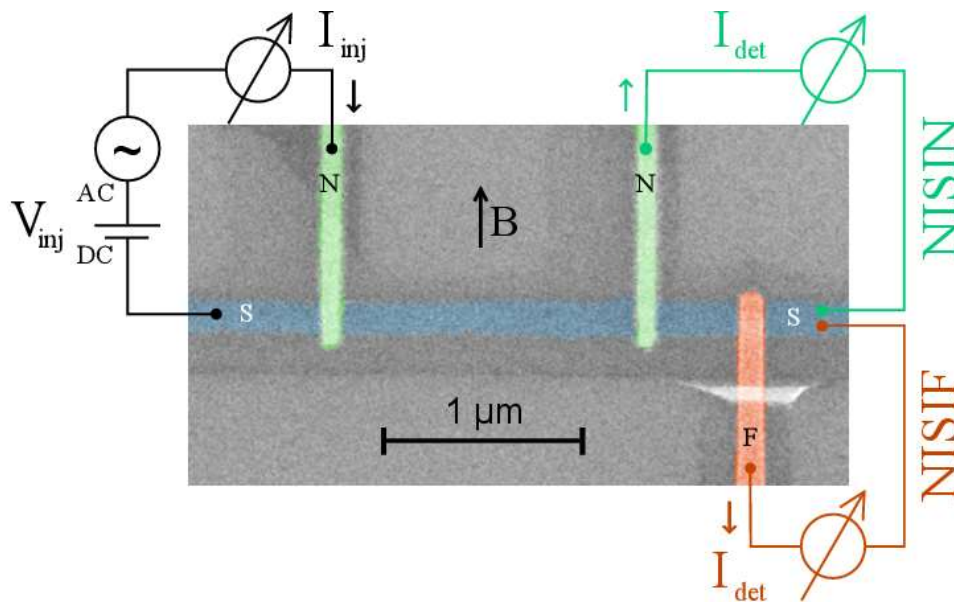
# Experiment – distance dependence



- peak area decreases with contact distance
- signal persists up to  $8 \mu m$
- relaxation length  $\lambda_S = 5 - 10 \mu m$  (compare to  $\lambda_N = 370 nm$ )
- what is the relaxation mechanism?



# Experiment – separating spin and charge

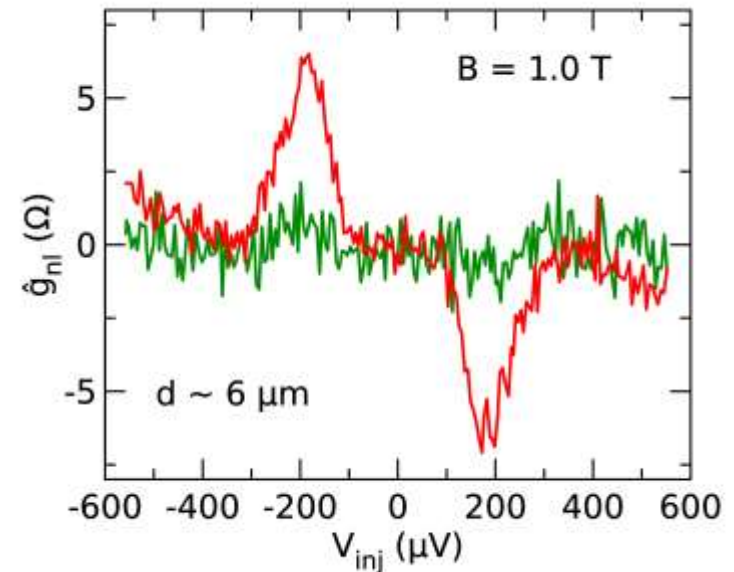


contact distance  $6 \mu\text{m}$

no more charge at all

but: still spin

- direct comparison of charge and spin signal in same sample
- spin injection from normal metal



Wolf et al., Phys. Rev. B **87**, 024517 (2013)

# Conclusions & Outlook

## Conclusions

- Long range spin transport in Zeeman-split superconductors
- model for relaxation needed

## Outlook

- thermoelectric effects
- manipulate and utilize spin currents

Thank you for your attention!

Hübler et al., Phys. Rev. B **81**, 184524 (2010)  
Hübler et al., Phys. Rev. Lett. **109**, 207001 (2012)  
Wolf et al., Phys. Rev. B **87**, 024517 (2013)



<http://www.int.kit.edu/english/606.php>